

## 6 Representation theory of the special unitary group $SU(N)$

### 6.1 Schur-Weyl duality — an overview

The Schur-Weyl duality is a powerful tool in representation theory that allows one to put the irreducible representations of the general linear group  $GL(N)$  (c.f. [Definition 6.1](#)) a vector space  $V^{\otimes n}$  with  $\dim V = N$  into 1-to-1 correspondence to the irreducible representations of the group  $S_n$  on  $V^{\otimes n}$ . In particular, it turns out that  $V^{\otimes n}$  decomposes as

$$V^{\otimes n} = \bigoplus_{\lambda} V_{\lambda} \otimes S_{\lambda}, \quad \lambda \vdash n, \quad (6.1)$$

where  $V_{\lambda}$  are the irreducible submodules of  $GL(N)$  on  $V^{\otimes n}$ , and  $S_{\lambda}$  are the so-called *Specht modules*, which describe the irreducible representations of  $S_n$  (c.f., e.g., [4] and other standard textbooks). The underlying reason for this is that the actions of  $GL(N)$  and  $S_n$  on  $V^{\otimes n}$  commute and, even more, the elements of  $GL(N)$  are a *complete set* of actions that commute with those of  $S_n$  and vice

In these lectures, we will go through the main points of the Schur-Weyl duality, paying particular attention to the role the Young projection operators play in the representation theory of  $GL(N)$ .

- We will begin by defining the general linear group  $GL(N)$  in [section 6.2](#).
- We will then define what we mean by an *invariant* of  $GL(N)$ , and show that these invariants are given by the elements of  $S_n$ , [section 6.3](#). In fact,  $S_n$  spans the algebra of invariants of  $GL(N)$ . An important ingredient to seeing this is the *double commutant theorem*, c.f. [section 6.3.1](#).
- Let  $v \in V^{\otimes n}$  be arbitrary. We will show that, for every  $\Theta \in \mathcal{Y}_n$ , the subspace

$$Y_{\Theta}v \quad (6.2)$$

is invariant and irreducible under the action of  $S_n$ . Hence,  $Y_{\Theta}v$  is an irreducible  $\mathbb{C}[S_n]$ -submodule and therefore corresponds to an irreducible representation of  $S_n$  on  $V^{\otimes n}$ , c.f. [section 6.4](#).

- Thereafter, we will show that, for every  $\Theta \in \mathcal{Y}_n$ , the subspace

$$Y_{\Theta}V^{\otimes n} \quad (6.3)$$

is invariant and irreducible under the action of  $GL(N)$  — the main ingredient to showing this is the fact that  $S_n$  spans the algebra of invariants of  $GL(N)$  on  $V^{\otimes n}$ . This shows that the Young projection operators  $Y_{\Theta}$  generate the irreducible ideals (and hence the irreducible representations) of  $GL(N)$  on  $V^{\otimes n}$ , [section 6.5](#).

- Finally, in [section 6.6](#), we will argue that the irreducible representations of  $GL(N)$  on  $V^{\otimes n}$  are precisely those of the special unitary group  $SU(N)$  on  $V^{\otimes n}$ . In other words, the Young projection operators on  $V^{\otimes n}$  give rise to all the irreducible ideals of  $SU(N)$  on  $V^{\otimes n}$ .
- We will end with an example, constructing the irreducible representations of  $GL(N)$  (hence also  $SU(N)$ ) and  $S_n$  on  $V^{\otimes n}$  for  $n = N = 3$ , in [section 6.7](#).

## 6.2 Basic definitions

### ■ Definition 6.1 – General linear group $\mathbf{GL}(V)$ (or $\mathbf{GL}(N)$ ):

Let  $V$  be a vector space of dimension  $N$  ( $N$  not necessarily finite). Consider the subset of  $\text{End}(V)$  of all invertible endomorphisms of  $V$ . This set forms a group called the general linear group on  $V$ , and we denote this group by  $\mathbf{GL}(V)$  or sometimes only  $\mathbf{GL}(N)$  if the vector space  $V$  is clear and we want to make the dimension of  $V$  explicit.

### ■ Definition 6.2 – Defining/fundamental representation of $\mathbf{GL}(V)$ :

Let  $\mathbf{GL}(V)$  be the general linear group acting on a vector space  $V$ . We can define a representation  $\gamma$  as

$$\begin{aligned}\gamma : \mathbf{GL}(V) &\rightarrow \text{End}(V) \\ \gamma(g) &\mapsto g\end{aligned}\tag{6.6}$$

since there is a well-defined action of  $\mathbf{GL}(V)$  on  $V$ . The representation  $\gamma$  is referred to as the defining representation of  $\mathbf{GL}(V)$  and has dimension  $\dim V = N$ .

## 6.3 Invariants of $\mathbf{GL}(N)$

Through the defining representation  $\gamma$  of  $\mathbf{GL}(V)$  on  $V$ , we can define a representation of  $\mathbf{GL}(V)$  on  $V^{\otimes n}$  via

$$g(v_1 \otimes v_2 \otimes \cdots \otimes v_n) := \gamma(g)v_1 \otimes \gamma(g)v_2 \otimes \cdots \otimes \gamma(g)v_n = gv_1 \otimes gv_2 \otimes \cdots \otimes gv_n\tag{6.7a}$$

Let  $\mathcal{A}$  be the symmetric group  $S_n$ , and let  $\mathcal{B}$  be the general linear group  $\mathrm{GL}(N)$ . Both these groups have a well-defined action on the vector space  $V^{\otimes n}$  (with  $\dim V = N$ ), and therefore both  $S_n$  and  $\mathrm{GL}(N)$  are subgroups of  $\mathrm{End}(V^{\otimes n})$ . Furthermore, we have seen in the beginning of [section 6.3](#) that the actions of  $S_n$  and  $\mathrm{GL}(N)$  commute on  $V^{\otimes n}$ , such that

$$\begin{aligned} \mathcal{A} = S_n \subset (\mathrm{GL}(N))' = \mathcal{B}' & \quad \text{and} \quad \mathcal{B} = \mathrm{GL}(N) \subset S_n' = \mathcal{A}' \\ \Rightarrow \quad \mathcal{A} \subset \mathcal{B}' & \quad \text{and} \quad \mathcal{B} \subset \mathcal{A}' \end{aligned} \quad (6.18)$$

It follows directly from Maschke's [Theorem 3.1](#) that  $V^{\otimes n}$  is a *completely reducible*  $S_n$ -module. That  $V^{\otimes n}$  is also a completely reducible  $\mathrm{GL}(N)$ -module follows from the *Peter-Weyl Theorem* [[16](#)] (which we will state without proof):

**Note 6.1: Peter-Weyl Theorem**

As we will explain in [Note 6.2](#),  $\mathrm{GL}(N)$  is a Lie group, which means that, in particular, it is a differentiable manifold. A manifold is said to be *compact* if it is compact as a topological space, that is every open cover has a finite subcover, *c.f.*, e.g., [[17](#)]. (*Very* loosely speaking, you may think of a cover of a manifold  $M$  as another manifold  $M'$  enclosing it. The requirement that every subcover is finite can be thought of that every submanifold  $N' \subset M'$  that is also a cover for  $M$  is finite, implying that  $M$  was finite to start off with. As an example, the unit sphere is a compact manifold, but an infinite plane would not be.) It turns out that the representations of such compact Lie groups are completely reducible:

■ **Theorem 6.2 – Reducible carrier spaces of compact groups (Peter-Weyl [[16](#)]):**  
*Let  $\varphi$  be a unitary representation of a compact group  $G$  on a complex Hilbert space  $H$ . Then  $H$  splits into an orthogonal direct sum of irreducible finite-dimensional unitary representations of  $G$ .*

Now, the group  $\mathrm{GL}(N)$  is not compact (as a manifold). However, as we will see in the later [section 6.6](#), its subgroup  $\mathrm{SU}(N)$  is compact. Furthermore, as we will argue in [section 6.6](#), the irreducible representations of  $\mathrm{GL}(N)$  are precisely those of  $\mathrm{SU}(N)$  and vice versa (*c.f.* [Theorem 6.5](#)) by means of the so-called *unitarian trick*, it follows that the Peter-Weyl [Theorem 6.2](#) does indeed apply to the group  $\mathrm{GL}(N)$  as well.

If all the things said in this note do not quite make sense to you yet, try re-reading this note *after* you have read [section 6.6](#) — this should clear things up for you.

Therefore  $V^{\otimes n}$  is completely reducible as an  $S_n$ -module, as well as as a  $\mathrm{GL}(N)$ -module. Therefore, the Double commutant [theorem 6.1](#) asserts that

$$\begin{aligned} \mathcal{A} = S_n = \mathcal{A}'' & \quad \text{and} \quad \mathcal{B} = \mathrm{GL}(N) = \mathcal{B}'' \\ \Rightarrow \quad \mathcal{A} = \mathcal{A}'' & \quad \text{and} \quad \mathcal{B} = \mathcal{B}'' \end{aligned} \quad (6.19)$$

## 6.6 The unitarian trick: irreducible representations of $SU(N)$ from $GL(N)$

Having established the connection between the irreducible representations of  $S_n$  and the irreducible representations of  $GL(N)$  on  $V^{\otimes n}$ ,

### ■ Definition 6.4 – Special unitary group $SU(N)$ :

Let  $GL(N)$  be the general linear group on a vector space  $V$  with  $\dim(V) = N$ . We define  $SU(N)$  to be the subset of matrices in  $GL(N)$  that are unitary (with respect to the canonical scalar product on  $V$ ) and have determinant 1,

$$SU(N) = \left\{ U \in GL(N) \mid UU^\dagger = \mathbb{1} \text{ and } \det U = 1 \right\} . \quad (6.41)$$

It can be shown that  $SU(N)$  is in fact a group (c.f. ) and we call it the special unitary group on  $V$ .

(The term *special* refers to the fact that the elements of  $SU(N)$  are unimodular, i.e. have determinant 1, and *unitary* refers to the property that  $UU^\dagger = \mathbb{1}$  for every  $U \in SU(N)$ .)

The following intermediate result will turn out to be quite useful when establishing the correspondence between the irreducible representations of  $GL(N)$  and the irreducible representations of  $SU(N)$  on  $V^{\otimes n}$ :

### ■ Proposition 6.3 – Square matrix decomposition:

Let  $M^{n \times n}$  be the space of all  $n \times n$  matrices with entries in  $\mathbb{C}$ , and let  $H^{n \times n}$  and  $A^{n \times n}$  be the spaces of all Hermitian, respectively, anti-Hermitian  $n \times n$  matrices with entries in  $\mathbb{C}$ . Then,

$$M^{n \times n} = H^{n \times n} \oplus A^{n \times n} . \quad (6.44)$$

*Proof of Proposition 6.3.* It is clear that  $H^{n \times n} + A^{n \times n} \subset M^{n \times n}$  since a sum of a Hermitian and an anti-Hermitian  $n \times n$  matrix will still yield an  $n \times n$  matrix. Conversely, let  $m \in M^{n \times n}$ . Then, we can write

$$m = \underbrace{\frac{1}{2}(m + m^\dagger)}_{=:m_+} + \underbrace{\frac{1}{2}(m - m^\dagger)}_{=:m_-} , \quad (6.45)$$

where  $m^\dagger$  is the Hermitian conjugate of  $m$ . Notice that  $m_+$  is Hermitian and  $m_-$  is anti-Hermitian. Thus, we also have that  $M^{n \times n} \subset H^{n \times n} + A^{n \times n}$ , implying that

$$M^{n \times n} = H^{n \times n} + A^{n \times n} \quad (6.46)$$

It remains to show that  $H^{n \times n} \cap A^{n \times n} = \{0\}$  to obtain the desired result: Let  $m \in H^{n \times n} \cap A^{n \times n}$ . Since  $m \in H^{n \times n}$ ,  $m$  is Hermitian. Furthermore, since also  $m \in A^{n \times n}$ ,  $m$  is anti-Hermitian. Therefore,

$$m \stackrel{m \in H^{n \times n}}{=} m^\dagger \stackrel{m \in A^{n \times n}}{=} -m \quad \Rightarrow \quad m = -m. \quad (6.47)$$

However, since all entries in  $m$  are elements of  $\mathbb{C}$ ,  $m = -m$  only holds for  $m$  being the zero matrix. Therefore,  $H^{n \times n} \cap A^{n \times n} = \{0\}$ , yielding

$$M^{n \times n} = H^{n \times n} \oplus A^{n \times n}, \quad (6.48)$$

as required. □

### Note 6.2: Generators of a Lie group

Without explicitly saying it, both  $\text{GL}(N)$  and  $\text{SU}(N)$  are *Lie groups*, which is to say that they are differentiable manifolds that also have a group structure defined on them. As differentiable manifolds, they have a tangent space at the identity called the *Lie algebra*. (The Lie algebras of the general linear group  $\text{GL}(N)$  and the special unitary group  $\text{SU}(N)$  are denoted as  $\mathfrak{gl}(N)$  and  $\mathfrak{su}(N)$ , respectively.) One can define a map, called the *exponential map* from the Lie algebra  $\mathfrak{g}$  to the Lie group  $\mathbf{G}$  as

$$\begin{aligned} \exp : \mathfrak{g} &\rightarrow \mathbf{G} \\ X &\mapsto e^{iX} \end{aligned} \quad (6.49)$$

(the  $i$  in the exponent is physics convention). It can be shown that, for certain Lie groups (including  $\text{GL}(N)$  and  $\text{SU}(N)$ , but not  $\text{SL}(N)$ ) this exponential map is *surjective*, implying that every element of the Lie group can be written as an exponential of the corresponding element in the Lie algebra! The Lie algebra is a linear space, and in particular it has a basis  $\{T_1, T_2, \dots, T_{\dim(\mathfrak{g})}\}$  such that every  $X \in \mathfrak{g}$  can be written as a direct sum of these basis vectors. Thus, each element  $g$  of the Lie group  $\mathbf{G}$  can be written as an exponential of a weighted sum of the basis vectors  $T_1, T_2, \dots, T_{\dim(\mathfrak{g})}$ ,

$$g = e^{i \sum_j \omega_j T_j}, \quad (6.50)$$

where the  $\omega_j$  are scalars. Hence, the Lie group  $\mathbf{G}$  is *generated* by the elements in  $\{T_1, T_2, \dots, T_{\dim(\mathfrak{g})}\}$  and we call the  $T_1, T_2, \dots, T_{\dim(\mathfrak{g})}$  the *generators of the Lie group*  $\mathbf{G}$ .

For this course, we will not discuss any differential geometry, we will merely use the fact that any element  $G$  of a Lie group can be written as an exponential  $G = e^{i \sum_j \omega_j T_j}$  where the  $T_j$  are the generators of the group. Readers, however, are encouraged to explore the topic more on their own, for example by reading through [17].

Now, notice that both  $\text{GL}(N)$  and  $\text{SU}(N)$  are matrix groups, and hence the exponential map is given by matrix exponentials. In other words, the generators of both  $\text{GL}(N)$  and  $\text{SU}(N)$  are square matrices.

By definition of  $\mathrm{SU}(N)$ , each  $U \in \mathrm{SU}(N)$  is unitary,  $UU^\dagger = \mathbb{1}$ . Writing  $U = e^{iX}$  for some  $X \in \mathfrak{su}(N)$ , we thus have that

$$\mathbb{1} \stackrel{!}{=} UU^\dagger = e^{iX} (e^{iX})^\dagger = e^{iX} e^{-iX^\dagger} = e^{i(X-X^\dagger)} \quad \Rightarrow \quad X - X^\dagger = 0. \quad (6.51)$$

Hence,  $X$  is Hermitian, and therefore the generators of  $\mathrm{SU}(N)$  must all be Hermitian. In fact, eq. (6.51) shows that any Hermitian matrix, when exponentiated, produces an element of  $\mathrm{SU}(N)$ . Thus, the subset of generators of  $\mathrm{GL}(N)$  that is Hermitian are the generators of  $\mathrm{SU}(N)$ .

As we have seen in [Proposition 6.3](#), every square matrix can be written as a sum of a Hermitian and an anti-Hermitian matrix. Notice that any Hermitian matrix can be made anti-Hermitian by multiplying it with  $i$  (and, obviously, any anti-Hermitian matrix with complex entries can be written as  $i$  times a Hermitian matrix). Therefore, the generators of  $\mathrm{GL}(N)$  are precisely the generators of  $\mathrm{SU}(N)$  plus  $i$  times the generators of  $\mathrm{SU}(N)$

Therefore, if  $\varphi : \mathrm{GL}(N) \rightarrow \mathrm{End}(W)$  is an irreducible representation of  $\mathrm{GL}(N)$  on  $W \subset V^{\otimes n}$  (with  $\dim V = N \geq n$ ), and we restrict this representation to the subgroup  $\mathrm{SU}(N) \subset \mathrm{GL}(N)$ , we merely only act  $\varphi$  on half of the generators of  $\mathrm{GL}(N)$ .<sup>6</sup> This clearly yields a representation of  $\mathrm{SU}(N)$

Suppose now that the resctriction of  $\varphi$  on  $\mathrm{SU}(N)$ ,  $\varphi|_{\mathrm{SU}(N)} : \mathrm{SU}(N) \rightarrow \mathrm{End}(W)$  is *not* irreducible as a representation of  $\mathrm{SU}(N)$ . Then one may decompose the carrier space  $W$  as a direct sum  $W = W_1 \oplus W_2$ , where each  $W_i$  is irreducible under the action of  $\mathrm{SU}(N)$ . Hence, any  $U \in \mathrm{SU}(N)$  can be written in the form

$$U := \begin{pmatrix} U_{W_1} & 0 \\ 0 & U_{W_2} \end{pmatrix}, \quad (6.52)$$

where  $U_{W_i}$  is of size  $\dim(W_i) \times \dim(W_i)$ , leaves the space  $W = W_1 \oplus W_2$  invariant. Since this is true for every  $U \in \mathrm{SU}(N)$ , the generators  $T_1, T_2, \dots, T_{\dim(\mathfrak{su}(N))}$  must block-diagonalize the matrices of  $\mathrm{SU}(N)$  on the space  $W$ . However, since we just reasoned that the  $T_j$  form a complete set of generators of  $\mathrm{GL}(N)$ , it follows that every  $g \in \mathrm{GL}(N)$  can be block-diagonalized on  $W$ , yielding  $W$  to be reducible. This is a contradiction as  $\varphi : \mathrm{GL}(N) \rightarrow W$  is known to be irreducible, and hence the restriction  $\varphi|_{\mathrm{SU}(N)} : \mathrm{SU}(N) \rightarrow \mathrm{End}(W)$  must be an irreducible representation of  $\mathrm{SU}(N)$  as well.

On the other hand, suppose  $\tilde{\varphi} : \mathrm{SU}(N) \rightarrow \mathrm{End}(\tilde{W})$  is an irreducible representation of  $\mathrm{SU}(N)$  on  $\tilde{W} \subset V^{\otimes n}$ . This representation can be extended to a representation of  $\mathrm{GL}(N)$ ,  $\tilde{\varphi}|^{\mathrm{GL}(N)} : \mathrm{GL}(N) \rightarrow \mathrm{End}(\tilde{W})$ , by including the action on  $i$  times the generators. Going through a similar chain of arguments as in the previous paragraph, it can be shown that  $\tilde{\varphi}|^{\mathrm{GL}(N)}$  is an irreducible representation of  $\mathrm{GL}(N)$ . Therefore, we have:

■ **Theorem 6.5 – Irreducible representations of  $\mathrm{SU}(N)$  and  $\mathrm{GL}(N)$ :**

*Any irreducible representation of  $\mathrm{GL}(N)$  is an irreducible representation of  $\mathrm{SU}(N)$  and vice versa.*



**Important:** Note that [Theorem 6.5](#) is very particular for the groups  $\mathrm{GL}(N)$  and  $\mathrm{SU}(N)$ . If one were to restrict an irreducible representation of  $\mathrm{GL}(N)$  to either the *orthogonal group*  $\mathrm{O}(N)$  or the *special orthogonal group*  $\mathrm{SO}(N)$  (both are subgroups of  $\mathrm{GL}(N)$ ), then the restricted representation is, in general, no longer irreducible. The reason for this is that both  $\mathrm{O}(N)$  and  $\mathrm{SO}(N)$  have additional invariants to  $\mathrm{GL}(N)$  on  $V^{\otimes n}$ , therefore

**Note 6.3: Invariants of  $SU(N)$**

The discussion thus far not only shows that the irreducible representations of  $SU(N)$  are precisely those of  $GL(N)$ , but further implies that also the linear invariants of  $SU(N)$  are those of  $GL(N)$ ,

$$\text{API}(GL(N), V^{\otimes n}) = \text{API}(SU(N), V^{\otimes n}) = \mathbb{C}[S_n] . \quad (6.53)$$

Furthermore, the Young projection operators  $Y_\Theta$  corresponding to the Young tableaux  $\Theta \in \mathcal{Y}_n$  generate the irreducible representations of  $SU(N)$  on  $V^{\otimes n}$ .

### 6.7 Example: The irreducible representations of $S_3$ and $SU(3)$ on $V^{\otimes 3}$

Consider a 3-dimensional vector space  $V$  with basis  $\{v_1, v_2, v_3\}$ . Forming the tensor product space  $V^{\otimes 3}$ , the basis of  $V$  induces a basis on  $V^{\otimes 3}$ , where each basis vector of  $V^{\otimes 3}$  is of the form

$$v_i \otimes v_j \otimes v_k \quad \text{for } i, j, k \in \{1, 2, 3\} ; \quad (6.54a)$$

clearly, this basis has size  $3^3 = 27$ . (In general, if  $\dim(V) = N$ , the tensor product space  $V^{\otimes n}$  has dimension  $N^n$ .) Introducing the shorthand notation

$$|ijk\rangle := v_i \otimes v_j \otimes v_k , \quad (6.54b)$$

the basis vectors of  $V^{\otimes 3}$  are given by

$$\begin{aligned} &|111\rangle , \quad |112\rangle , \quad |121\rangle , \quad |211\rangle , \quad |122\rangle , \quad |221\rangle , \quad |212\rangle , \\ &|222\rangle , \quad |113\rangle , \quad |131\rangle , \quad |311\rangle , \quad |133\rangle , \quad |331\rangle , \quad |313\rangle , \\ &|333\rangle , \quad |223\rangle , \quad |232\rangle , \quad |322\rangle , \quad |233\rangle , \quad |332\rangle , \quad |323\rangle , \\ &|123\rangle , \quad |132\rangle , \quad |213\rangle , \quad |231\rangle , \quad |312\rangle , \quad |321\rangle . \end{aligned} \quad (6.55)$$

As usual, the action of a permutation  $\rho$  in  $S_3$  on a (basis) vector in  $V^{\otimes 3}$  is realized through the permutation of its tensor indices, for example

$$(123) |123\rangle = \frac{1}{2} = \frac{3}{2} = |312\rangle , \quad (6.56a)$$

and the action of any  $U \in SU(3)$  on  $|ijk\rangle$  yields

$$U |ijk\rangle = |U(i)U(j)U(k)\rangle . \quad (6.56b)$$

Let us now study the irreducible representations of  $S_3$  and  $SU(3)$  on  $V^{\otimes 3}$ :

As we have seen in this section, the irreducible representations of  $S_n$  and  $SU(N)$  on  $V^{\otimes n}$  are generated by the Young projection operators corresponding to the Young tableaux in  $\mathcal{Y}_n$ . Here, for  $n = 3$ , we have that

$$\mathcal{Y}_3 = \left\{ \begin{array}{|c|c|c|} \hline 1 & 2 & 3 \\ \hline \end{array} , \quad \begin{array}{|c|c|} \hline 1 & 2 \\ \hline 3 \\ \hline \end{array} , \quad \begin{array}{|c|c|} \hline 1 & 3 \\ \hline 2 \\ \hline \end{array} , \quad \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline \end{array} \right\} . \quad (6.57)$$

First, let us consider the two Young tableaux of shape  $\begin{array}{|c|c|} \hline & \\ \hline & \\ \hline \end{array}$ ,

$$\Theta = \begin{array}{|c|c|} \hline 1 & 2 \\ \hline 3 & \\ \hline \end{array} \quad \text{and} \quad \Phi = \begin{array}{|c|c|} \hline 1 & 3 \\ \hline 2 & \\ \hline \end{array} . \quad (6.58)$$

According to [Theorem 5.2](#), the Young projection operators  $Y_\Theta$  and  $Y_\Phi$  produce an equivalent 2-dimensional irreducible representations of  $S_3$  on  $V^{\otimes 3}$ . Let us see this explicitly:

First, we need to consider the action of both  $Y_\Theta$  and  $Y_\Phi$  on the basis vectors (6.55) of  $V^{\otimes 3}$ . For the Young projection operator corresponding to the tableau  $\Theta$

$$Y_\Theta = \frac{4}{3} \begin{array}{c} \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \end{array} = \frac{1}{3} \left( \quad - \quad + \quad - \quad \right) , \quad (6.59)$$

notice that, for  $(13) \in S_3$ ,

$$Y_\Theta(13) = \frac{4}{3} \begin{array}{c} \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \end{array} = \frac{4}{3} \begin{array}{c} \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \end{array} = -\frac{4}{3} \begin{array}{c} \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \end{array} = -Y_\Theta , \quad (6.60a)$$

such that, for each vector  $|ijk\rangle$ , we have that

$$Y_\Theta |ijk\rangle = -Y_\Theta(13) |ijk\rangle = -Y_\Theta |kji\rangle , \quad \text{and hence} \quad Y_\Theta |iji\rangle = 0 . \quad (6.60b)$$

Therefore, acting  $Y_\Theta$  on the basis vectors of  $V^{\otimes 3}$  given in eq. (6.55) yields the following 8 *nonzero, linearly independent* vectors:

$$Y_\Theta |112\rangle = -Y_\Theta |211\rangle = \frac{1}{3} (2|112\rangle - |211\rangle - |121\rangle) \quad (6.61a)$$

$$Y_\Theta |113\rangle = -Y_\Theta |311\rangle = \frac{1}{3} (2|113\rangle - |311\rangle - |131\rangle) \quad (6.61b)$$

$$Y_\Theta |223\rangle = -Y_\Theta |322\rangle = \frac{1}{3} (2|223\rangle - |322\rangle - |232\rangle) \quad (6.61c)$$

$$Y_\Theta |221\rangle = -Y_\Theta |122\rangle = \frac{1}{3} (2|221\rangle - |122\rangle - |212\rangle) \quad (6.61d)$$

$$Y_\Theta |331\rangle = -Y_\Theta |133\rangle = \frac{1}{3} (2|331\rangle - |133\rangle - |313\rangle) \quad (6.61e)$$

$$Y_\Theta |332\rangle = -Y_\Theta |233\rangle = \frac{1}{3} (2|332\rangle - |233\rangle - |323\rangle) \quad (6.61f)$$

$$Y_\Theta |123\rangle = -Y_\Theta |321\rangle = \frac{1}{3} (|123\rangle - |321\rangle + |213\rangle - |231\rangle) \quad (6.61g)$$

$$Y_\Theta |132\rangle = -Y_\Theta |231\rangle = \frac{1}{3} (|132\rangle - |231\rangle + |312\rangle - |321\rangle) , \quad (6.61h)$$

where

$$Y_\Theta |213\rangle = -Y_\Theta |312\rangle = Y_\Theta |123\rangle - Y_\Theta |132\rangle . \quad (6.62)$$

Similarly, for the Young projection operator  $Y_\Phi$ , we notice the following symmetry

$$Y_\Phi(12) = \frac{4}{3} \begin{array}{c} \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \end{array} = -\frac{4}{3} \begin{array}{c} \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \end{array} = -Y_\Phi , \quad (6.63a)$$

such that, for each vector  $|ijk\rangle$ , we have that

$$Y_\Phi |ijk\rangle = -Y_\Phi(12) |ijk\rangle = -Y_\Phi |jik\rangle , \quad \text{and hence} \quad Y_\Phi |iij\rangle = 0 . \quad (6.63b)$$



Hence,

$$Y_{\Phi} |121\rangle = -Y_{\Phi} |211\rangle = \frac{1}{3} (2 |121\rangle - |211\rangle - |112\rangle) \quad (6.64a)$$

$$Y_{\Phi} |131\rangle = -Y_{\Phi} |311\rangle = \frac{1}{3} (2 |131\rangle - |311\rangle - |113\rangle) \quad (6.64b)$$

$$Y_{\Phi} |232\rangle = -Y_{\Phi} |322\rangle = \frac{1}{3} (2 |232\rangle - |322\rangle - |223\rangle) \quad (6.64c)$$

$$Y_{\Phi} |212\rangle = -Y_{\Phi} |122\rangle = \frac{1}{3} (2 |212\rangle - |122\rangle - |221\rangle) \quad (6.64d)$$

$$Y_{\Phi} |313\rangle = -Y_{\Phi} |133\rangle = \frac{1}{3} (2 |313\rangle - |133\rangle - |331\rangle) \quad (6.64e)$$

$$Y_{\Phi} |323\rangle = -Y_{\Phi} |233\rangle = \frac{1}{3} (2 |323\rangle - |233\rangle - |332\rangle) \quad (6.64f)$$

$$Y_{\Phi} |123\rangle = -Y_{\Phi} |213\rangle = \frac{1}{3} (|123\rangle - |213\rangle + |321\rangle - |312\rangle) \quad (6.64g)$$

$$Y_{\Phi} |132\rangle = -Y_{\Phi} |312\rangle = \frac{1}{3} (|132\rangle - |312\rangle + |231\rangle - |213\rangle) , \quad (6.64h)$$

and again

$$Y_{\Phi} |231\rangle = -Y_{\Phi} |321\rangle = -Y_{\Phi} |123\rangle - Y_{\Phi} |132\rangle . \quad (6.65)$$

Consider now the irreducible representation of  $S_3$  generated by  $Y_{\Theta}$ . If we define an operator  $\mathcal{T}_{\Phi\Theta}$  as

$$\mathcal{T}_{\Phi\Theta} := (23)Y_{\Theta} = \text{diagram} , \quad (6.66)$$

then we have that, for all  $\rho \in S_3$ ,

$$\rho Y_{\Theta} = c_1 Y_{\Theta} + c_2 \mathcal{T}_{\Phi\Theta} , \quad \text{where } c_1, c_2 \in \mathbb{C} , \quad (6.67)$$

which is to say that  $Y_{\Theta}$  generates a 2-dimensional submodule of  $V^{\otimes 3}$ . Explicitly:

$$\text{id}_3 Y_{\Theta} = \frac{4}{3} \cdot \text{diagram} = \frac{4}{3} \text{diagram} = Y_{\Theta} \quad (6.68a)$$

$$(12)Y_{\Theta} = \frac{4}{3} \cdot \text{diagram} = \frac{4}{3} \text{diagram} = Y_{\Theta} \quad (6.68b)$$

$$(23)Y_{\Theta} = \frac{4}{3} \cdot \text{diagram} = \frac{4}{3} \text{diagram} = \mathcal{T}_{\Phi\Theta} \quad (6.68c)$$

$$(13)Y_{\Theta} = \frac{4}{3} \cdot \text{diagram} = -\frac{4}{3} \text{diagram} - \frac{4}{3} \text{diagram} = -Y_{\Theta} - \mathcal{T}_{\Phi\Theta} \quad (6.68d)$$

$$(123)Y_{\Theta} = \frac{4}{3} \cdot \text{diagram} = -\frac{4}{3} \text{diagram} - \frac{4}{3} \text{diagram} = -Y_{\Theta} - \mathcal{T}_{\Phi\Theta} \quad (6.68e)$$

$$(132)Y_{\Theta} = \frac{4}{3} \cdot \text{diagram} = \frac{4}{3} \text{diagram} = \mathcal{T}_{\Phi\Theta} ; \quad (6.68f)$$

eq. (6.68d) can be easiest seen as follows:

$$\begin{aligned}
(13)Y_\Theta &= \frac{4}{3} \cdot \text{diagram} = -\frac{4}{3} \cdot \text{diagram} = -\frac{4}{3} \cdot \text{diagram} \\
&= -\frac{4}{3} \cdot \text{diagram} = -2\frac{1}{2}\frac{4}{3} \left( \text{diagram} + \underbrace{\text{diagram}}_{=0} \right) = -2\frac{4}{3} \cdot \text{diagram} \\
&= -2\frac{1}{2}\frac{4}{3} \left( \text{diagram} + \text{diagram} \right) = -Y_\Theta - \mathcal{T}_{\Phi\Theta}, \tag{6.69}
\end{aligned}$$

and eq. (6.68e) follows from

$$(123)Y_\Theta = \frac{4}{3} \cdot \text{diagram} = \frac{4}{3} \cdot \text{diagram} = \frac{4}{3} \cdot \text{diagram} = (13)Y_\Theta. \tag{6.70}$$

Hence,  $Y_\Theta$  indeed generates a 2-dimensional  $S_3$ -submodule  $Y_\Theta \mathbf{v}$  of  $V^{\otimes 3}$ , as claimed in eq. (6.67). Since there exist 8 linearly independent vectors  $Y_\Theta \mathbf{v}$ , this representation is contained 8 times within the regular representation of  $S_3$ , and we say that *the 2-dimensional  $S_3$ -module generated by  $Y_\Theta$  has multiplicity 8.*

We will show that the tableau  $\Phi$  given in eq. (6.58) also gives rise to a 2-dimensional  $S_3$  module with multiplicity 8, which will be shown to be completely equivalent to that generated by  $Y_\Theta$ : The Young projection operator  $Y_\Phi$  is given by

$$Y_\Phi = \frac{4}{3} \cdot \text{diagram} = \frac{1}{3} \left( \text{diagram} + \text{diagram} - \text{diagram} \right). \tag{6.71}$$

Defining

$$\mathcal{T}_{\Phi\Theta} := (23)Y_\Theta = \frac{4}{3} \cdot \text{diagram}, \tag{6.72}$$

one may again show that  $Y_\Phi$  generates a 2-dimensional submodule of  $V^{\otimes 3}$  as

$$\rho Y_\Phi = k_1 Y_\Phi + k_2 \mathcal{T}_{\Phi\Theta}, \quad \text{for every } \rho \in S_3, \text{ where } c_1, c_2 \in \mathbb{C}. \tag{6.73}$$

In particular,

$$\text{id}_3 Y_\Theta = \frac{4}{3} \cdot \text{diagram} = \frac{4}{3} \cdot \text{diagram} = Y_\Theta \tag{6.74a}$$

$$(12)Y_\Theta = \frac{4}{3} \cdot \text{diagram} = -\frac{4}{3} \cdot \text{diagram} - \frac{4}{3} \cdot \text{diagram} = -Y_\Theta - \mathcal{T}_{\Phi\Theta} \tag{6.74b}$$

$$(23)Y_\Theta = \frac{4}{3} \cdot \text{diagram} = \frac{4}{3} \cdot \text{diagram} = \mathcal{T}_{\Phi\Theta} \tag{6.74c}$$

$$(13)Y_\Theta = \frac{4}{3} \cdot \text{diagram} = \frac{4}{3} \cdot \text{diagram} = Y_\Theta \tag{6.74d}$$

$$(123)Y_\Theta = \frac{4}{3} \cdot \text{diagram} = \frac{4}{3} \cdot \text{diagram} = \mathcal{T}_{\Phi\Theta} \tag{6.74e}$$

$$(132)Y_\Theta = \frac{4}{3} \cdot \text{diagram} = -\frac{4}{3} \cdot \text{diagram} - \frac{4}{3} \cdot \text{diagram} = -Y_\Theta - \mathcal{T}_{\Phi\Theta}. \tag{6.74f}$$

Notice that this agrees with the dimension formula for the irreducible representations of  $GL(N)$  (hence, also  $SU(N)$ ) given in [section 6.5.1](#), as

$$\dim Y_{\Theta} = \text{tr} \left( \frac{4}{3} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \right) = \frac{4}{3} \frac{N(N^2 - 1)}{4} \stackrel{N=3}{=} 8 \quad (6.80a)$$

$$\dim Y_{\Phi} = \text{tr} \left( \frac{4}{3} \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \right) = \frac{4}{3} \frac{N(N^2 - 1)}{4} \stackrel{N=3}{=} 8 . \quad (6.80b)$$

One can perform a similar analysis for the remaining two tableaux in [\(6.57\)](#). For example, for

$$\Psi = \begin{array}{|c|c|c|} \hline 1 & 2 & 3 \\ \hline \end{array} , \quad (6.81)$$

the corresponding Young projection operator  $Y_{\Psi}$  is given by

$$Y_{\Psi} = \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \Rightarrow \text{tr} \left( \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \right) = \frac{N(N+1)(N+2)}{6} \stackrel{N=3}{=} 10 . \quad (6.82)$$

Since  $\Psi$  is the unique Young tableau of shape  $\square\square\square$ ,  $\Psi$  gives rise to 10 equivalent 1-dimensional irreducible representations of  $S_n$  on  $V^{\otimes 3}$ , and a unique 10-dimensional irreducible representation of  $SU(N)$  on  $V^{\otimes 3}$ .

Lastly, for the tableay

$$\Xi = \begin{array}{|c|} \hline 1 \\ \hline 2 \\ \hline 3 \\ \hline \end{array}, \quad (6.83)$$

the corresponding Young projection operator  $Y_\Xi$  is given by

$$Y_\Xi = \begin{array}{|c|} \hline \mathbf{1} \\ \hline \mathbf{2} \\ \hline \mathbf{3} \\ \hline \end{array} \Rightarrow \text{tr} \left( \begin{array}{|c|} \hline \mathbf{1} \\ \hline \mathbf{2} \\ \hline \mathbf{3} \\ \hline \end{array} \right) = \frac{N(N-1)(N-2)}{6} \stackrel{N=3}{=} 1. \quad (6.84)$$

Since  $\Xi$  again is the unique Young tableau of shape  $\begin{array}{|c|} \hline \square \\ \hline \square \\ \hline \square \\ \hline \end{array}$ ,  $\Xi$  gives rise one more 1-dimensional irreducible representations of  $S_n$ , which is *inequivalent* to the 1-dimensional representations obtained from  $\Psi$ ! Furthermore,  $Y_\Xi$  a unique 1-dimensional irreducible representation of  $SU(N)$  on  $V^{\otimes 3}$  — a 1-dimensional representation of  $SU(N)$  is also called a *singlet representation*.

Notice, however, that for  $\dim(V) = N \leq 2$ ,  $\text{tr}(Y_\Xi) = 0$ , and the tableau  $\Xi$  does not give rise to any representations of  $S_3$  or  $SU(N)$  on  $V^{\otimes 3}$ . In fact, for  $N \leq 2$ , the Young projection operator  $Y_\Xi$  becomes *dimensionally zero*, c.f. [Note 5.3](#).

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# The Special Unitary Group, Birdtracks, and Applications in QCD

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